

## Extreme Swelling of Lamellar Phases

In a recent Letter,<sup>1</sup> Larche *et al.* reported an unusual observation of extremely dilute lamellar phases of lyotropic liquid crystals. They evoke the steric repulsion between fluctuating lamellae<sup>2</sup> as a possible origin of stabilization of such phases. Recently we have studied<sup>3</sup> an analogous problem of the interaction between membranes, and have predicted the existence of *critical unbinding* transitions, driven by *fluctuations*, between a state in which the membranes are bound and the state in which they are separated. It is argued in this Comment that the process of swelling of a stack of lamellae can be viewed as a new kind of unbinding transition, named here *complete unbinding*. This unifying theoretical picture clarifies the notion and the use of the *steric repulsive potential*,<sup>2</sup> and shows that the detailed study of the swollen phases can bring new insight into the nature of molecular interactions in lyotropic liquid crystals.

Our theory is based on a simple effective Hamiltonian for a pair of lamellae,<sup>3</sup>

$$\mathcal{H}\{l(\mathbf{x})\} = \int d^{d-1}x \left[ \frac{1}{2} \kappa_0 (\nabla^2 l)^2 + V(l(\mathbf{x})) \right] / k_B T, \quad (1)$$

where  $l(\mathbf{x})$  is the separation between the lamellae and  $\kappa_0$  is the rigidity constant. To describe the swelling process one should consider *constrained systems*, since the mean separation  $\bar{l}$  between the lamellae is determined by the global composition of the mixture, or by the osmotic pressure inside the lamellar phase. We therefore add to the microscopic interaction  $V_0(l)$ , an extra pressure-like term,  $Pl$ . The swelling process can then be viewed as relaxing the constraint, i.e., as taking the limit  $P \rightarrow 0$  (see Fig. 1). If the molecular potential,  $V_0(l)$ , has a sufficiently small minimum, then the lamellae will separate completely in this limit. This happens, for example, when the Hamaker constant  $W$ , which governs the attractive part  $V_A(l)$ , is smaller than some critical value  $W_c$ . This new transition is characterized by  $\bar{l} \sim P^{-\psi}$  ( $P \rightarrow 0$ ). A self-consistent perturbation treatment of the model (1) which we have performed<sup>4</sup> yields the following results: (i) For sufficiently short-ranged molecular potentials  $V_0(l)$  [or more precisely for the potentials such that for large  $l$ ,  $V(l)l^{2(d-1)/(5-d)} \rightarrow 0$ ] the critical exponent  $\psi$  is given by  $\psi = (5-d)/(3+d)$ , which implies  $\psi = \frac{1}{3}$  in three dimensions; and (ii) if one includes in the potential  $V_0(l)$  an unscreened electrostatic repulsive interaction, then  $\psi = \frac{1}{2}$  ( $d=3$ ).

The result (i) can also be obtained in  $d=3$  by our simply adding the effective steric repulsion term,<sup>2</sup>  $V_{st} \sim (k_B T)^2 / \kappa l^2$ , to the potential  $V(l)$  and then using a simple minimization procedure. This approach is *not al-*

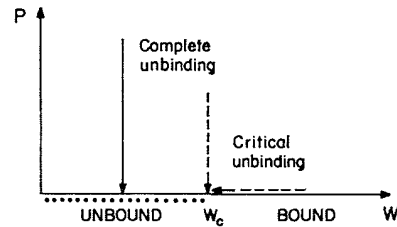


FIG. 1. Complete unbinding occurs for  $W < W_c$  as  $P \rightarrow 0$ ; the critical unbinding transition occurs at the point  $(W, P) = (W_c, 0)$ . Instead of the Hamaker constant  $W$ , one could vary the rigidity or the inverse temperature.

ways correct, however, as shown for the case of critical unbinding transitions (see Ref. 3), and one should in general perform a complete statistical (e.g., renormalization group) treatment of the model (1). On the other hand, our theory<sup>4</sup> shows that for *complete unbinding* the fluctuations are weak enough that their effect can be taken into account by this simple approach.

The model (1) can be generalized to the case of a *stack of lamellae*, and an effective Gaussian Hamiltonian can be derived.<sup>4</sup> This Hamiltonian leads to the prediction of *quasi long-range order* with characteristic exponent  $X_m$ . Asymptotically, for separations  $\bar{l}$  large compared to the thickness of the lamellae  $\delta$ , we predict that  $X_m \sim m^2$  is a pure number for case (i), and  $X_m \sim \bar{l}^{-1/2}$  for case (ii). Thus it should be possible to distinguish experimentally the situation where the “hyper-swollen” phases are stabilized by fluctuations of lamellae, or by unscreened electrostatic forces.<sup>5</sup>

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<sup>2</sup>W. Helfrich, *Z. Naturforsch. Teil A* **33**, 305 (1978).

<sup>3</sup>R. Lipowsky and S. Leibler, *Phys. Rev. Lett.* **56**, 2541 (1986).

<sup>4</sup>S. Leibler and R. Lipowsky, *Phys. Rev. B* (to be published).

<sup>5</sup>This has been confirmed in recent experiments, see D. Roux, in *Proceedings of Les Houches Conference on Amphiphilic Films*, Les Houches, France, February, 1987 (to be published).