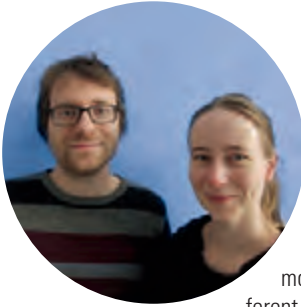


MOLECULAR MOTORS AND FILAMENTS

Cargo Transport by Teams of Molecular Motors: Elastic Coupling and Interference Regimes



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Molecular Motors: Network Representa-
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Active transport by molecular motors is ubiquitous in all eukaryotic cells. These motor proteins convert chemical energy into mechanical work which enables them to transport cargoes such as RNAs, protein complexes or organelles. Such intracellular transport is often driven cooperatively by several molecular motors, which may belong to the same or to different motor species like kinesin, dynein or myosin.

These motors step stochastically along filaments, from which they unbind after a finite run length. Understanding how these motors interact and how their movements are coordinated and regulated is a central and challenging problem of cellular transport [1]. To establish a general theoretical framework for elucidating and analysing such transport processes, we recently introduced two complementary approaches: (i) a detailed enzymatic chemo-mechanical description that starts from the nucleotide states of the single motors [2] and (ii) a coarse-grained description considering the single motor as a stochastic stepper [3,4]. Such theoretical approaches integrate the well-established properties of individual motors into a predictive theory for cooperative transport.

Network Representation of Motor Pairs

Cargo transport by two coupled molecular motors is studied using a chemomechanical network for the complete transport system and analyzing the trajectories generated by this network. The theoretical description starts from the different nucleotide states of a single motor supplemented by chemical and mechanical transitions between these states. As an instructive example, we focus on kinesin-1, for which a detailed chemomechanical network has been developed previously [5]. The motor pair system considered here consists of two kinesin motors, a leading and a trailing motor, which are attached to the same cargo and walk on the same filament. Each motor can unbind from and rebind to the filament individually. As a consequence, the cargo is actively pulled by either one or two motors. A mechanical step of one of the motors during a 2-motor run leads to an elastic interaction force between the two motors described by the extension of a linear spring. The state space of the elastically coupled motor pair is characterized by three variables, the chemical states i_{le} and i_{lr} of the leading and trailing motor and of the extension ΔL of the motor-motor separation [2]. The resulting network has a layer structure as shown in Fig. 1, where each layer corresponds to a constant value of ΔL . Any 1-motor run occurs on one of the boundary lines of the network and may be terminated either by unbinding of the active motor which

leads to the unbound motor pair state, or by the rebinding of the inactive motor and a subsequent 2-motor run. Mechanical steps during 2-motor runs lead to transitions between neighbouring ΔL -layers. Even though this motor pair network has a complex structure, it involves, apart from the single motor parameters, only two additional parameters, the coupling parameter K and the single motor rebinding rate π_{si} .

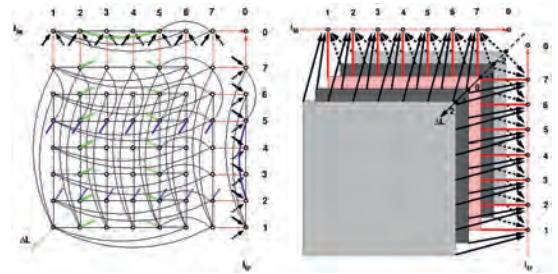


Fig. 1: State space of a motor pair described by the individual motor states i_{le} and i_{lr} and the extension ΔL of the motor-motor separation. (left) Detailed description of the layer with $\Delta L=0$ and (right) stack of five ΔL -layers.

Activity States and Motor Pair Parameters

In stochastic simulations, we studied the trajectories generated from the chemomechanical network of the motor pair as a function of the single motor rebinding rate π_{si} and the elastic coupling parameter K [2]. In experimental studies, the values of these two parameters are usually not known but have to be determined in a consistent manner. Our theory shows that one can determine these two parameters by measuring the average run times during 1- and 2-motor runs of cargo trajectories. Alternatively, individual motor trajectories and the properties of the ΔL -distribution can be used to deduce the two unknown parameters. Which activity state is dominant during a motor pair walk also depends on these two parameters. The corresponding activity diagram in Fig. 2 shows the crossover line which separates the parameter regime, in which 1-motor runs dominate the cargo run, from the regime, in which 2-motor runs are more likely. From individual motor trajectories, one can deduce the distribution of the extension ΔL of the motor-motor separation during 2-motor runs as shown in Fig. 2. Within the studied range for the coupling parameter, the number of accessible ΔL values varies by one order of magnitude. The maximal values of ΔL observed in the simulations determines the size of the network in Fig. 1.

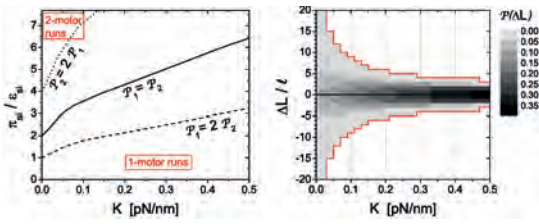


Fig. 2: (left) Activity regimes of a motor pair as a function of elastic coupling parameter K and single motor rebinding rate π_{on} . (right) Probability distribution $P(\Delta L)$ for the extension ΔL of the motor-motor separation as a function of the coupling parameter K .

Distinct Transport Regimes for Elastically Coupled Motors

The case of cargo transport by two identical motors involves an elastic coupling between the motors that can reduce the motors' velocity and/or the binding time to the filament. We show that this elastic coupling leads, in general, to four distinct transport regimes characterized by the motor pair's average velocity v_2 and its average binding time t_2 , during which the two motors remain simultaneously bound to the filament [3]. Both quantities depend on the single motor dynamics and on the strength K of their elastic coupling. Thus, strongly coupled and/or fast motors can quickly build up a large strain force that pulls one of the motors from the filament, while weakly coupled and/or slow motors may unbind spontaneously before reaching such a large force. The motor pair dynamics are governed by the interplay of three different forces: the stall force F_s , which corresponds to the maximal force that a single motor can generate, the detachment force F_d , which is the typical force that one motor can sustain for an extended period of time, and the scale F_K for the elastic strain forces between the two motors.

Using a continuous-time Markov process to describe the single motors as stochastic steppers, we calculate the average binding time t_2 and the velocity v_2 for two active motors and identify four different transport regimes, see Fig. 3. We estimate the crossover lines between these regimes from time scale arguments for the strain force generation which, in addition, allows us to obtain an intuitive understanding of the mutual motor-motor interference.

We apply our framework to predict the behavior of different pairs of molecular motors based on typical parameters from single motor experiments. In addition to a weak coupling regime, kinesin and dynein motors are found to exhibit a strong coupling and an enhanced unbinding regime, whereas myosin motors are predicted to attain a reduced velocity regime. All of these regimes can be explored experimentally by varying the elastic coupling parameter K . Our theory is consistent with the available experimental data for a kinesin-1 and myosin V.

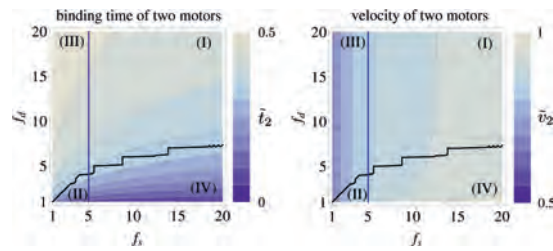


Fig. 3: (left) Average binding time and (right) velocity of a cargo pulled by two active motors in units of the single motor binding time and single motor velocity, respectively. Both quantities are displayed in a colored contour plot as functions of the rescaled single motor stall force $f_s = F_s/F_K$ and the rescaled detachment force $f_d = F_d/F_K$, where the force scale $F_K = K/2l$ is the elastic strain force between the motors per motor step with step size l . The solid blue line separates the region with and without a reduced velocity, whereas the solid black line separates the region of spontaneous and force-induced unbinding. These lines define four distinct transport regimes: (I) weak coupling with essentially unchanged velocity and binding time; (II) strong coupling with both reduced binding time and reduced velocity; (III) reduced velocity with no effect on the binding time and (IV) enhanced unbinding with no effect on the velocity.

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