

Extreme Swelling of Lamellar Phases

In a recent Letter,¹ Larche *et al.* reported an unusual observation of extremely dilute lamellar phases of lyotropic liquid crystals. They evoke the steric repulsion between fluctuating lamellae² as a possible origin of stabilization of such phases. Recently we have studied³ an analogous problem of the interaction between membranes, and have predicted the existence of *critical unbinding* transitions, driven by *fluctuations*, between a state in which the membranes are bound and the state in which they are separated. It is argued in this Comment that the process of swelling of a stack of lamellae can be viewed as a new kind of unbinding transition, named here *complete unbinding*. This unifying theoretical picture clarifies the notion and the use of the *steric repulsive potential*,² and shows that the detailed study of the swollen phases can bring new insight into the nature of molecular interactions in lyotropic liquid crystals.

Our theory is based on a simple effective Hamiltonian for a pair of lamellae,³

$$\mathcal{H}\{l(\mathbf{x})\} = \int d^{d-1}x \left[\frac{1}{2} \kappa_0 (\nabla^2 l)^2 + V(l(\mathbf{x})) \right] / k_B T, \quad (1)$$

where $l(\mathbf{x})$ is the separation between the lamellae and κ_0 is the rigidity constant. To describe the swelling process one should consider *constrained systems*, since the mean separation \bar{l} between the lamellae is determined by the global composition of the mixture, or by the osmotic pressure inside the lamellar phase. We therefore add to the microscopic interaction $V_0(l)$, an extra pressure-like term, Pl . The swelling process can then be viewed as relaxing the constraint, i.e., as taking the limit $P \rightarrow 0$ (see Fig. 1). If the molecular potential, $V_0(l)$, has a sufficiently small minimum, then the lamellae will separate completely in this limit. This happens, for example, when the Hamaker constant W , which governs the attractive part $V_A(l)$, is smaller than some critical value W_c . This new transition is characterized by $\bar{l} \sim P^{-\psi}$ ($P \rightarrow 0$). A self-consistent perturbation treatment of the model (1) which we have performed⁴ yields the following results: (i) For sufficiently short-ranged molecular potentials $V_0(l)$ [or more precisely for the potentials such that for large l , $V(l)l^{2(d-1)/(5-d)} \rightarrow 0$] the critical exponent ψ is given by $\psi = (5-d)/(3+d)$, which implies $\psi = \frac{1}{3}$ in three dimensions; and (ii) if one includes in the potential $V_0(l)$ an unscreened electrostatic repulsive interaction, then $\psi = \frac{1}{2}$ ($d=3$).

The result (i) can also be obtained in $d=3$ by our simply adding the effective steric repulsion term,² $V_{st} \sim (k_B T)^2 / \kappa l^2$, to the potential $V(l)$ and then using a simple minimization procedure. This approach is *not al-*

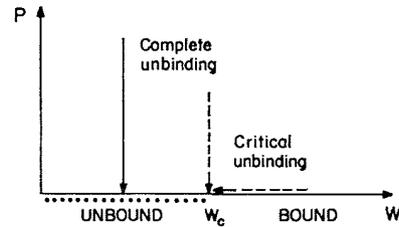


FIG. 1. Complete unbinding occurs for $W < W_c$ as $P \rightarrow 0$; the critical unbinding transition occurs at the point $(W, P) = (W_c, 0)$. Instead of the Hamaker constant W , one could vary the rigidity or the inverse temperature.

ways correct, however, as shown for the case of critical unbinding transitions (see Ref. 3), and one should in general perform a complete statistical (e.g., renormalization group) treatment of the model (1). On the other hand, our theory⁴ shows that for *complete unbinding* the fluctuations are weak enough that their effect can be taken into account by this simple approach.

The model (1) can be generalized to the case of a *stack of lamellae*, and an effective Gaussian Hamiltonian can be derived.⁴ This Hamiltonian leads to the prediction of *quasi long-range order* with characteristic exponent X_m . Asymptotically, for separations \bar{l} large compared to the thickness of the lamellae δ , we predict that $X_m \sim m^2$ is a pure number for case (i), and $X_m \sim \bar{l}^{-1/2}$ for case (ii). Thus it should be possible to distinguish experimentally the situation where the “hyper-swollen” phases are stabilized by fluctuations of lamellae, or by unscreened electrostatic forces.⁵

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¹F. C. Larche, J. Appell, G. Porte, P. Bassereau, and J. Maignan, Phys. Rev. Lett. **56**, 1700 (1986). See also C. R. Safinya, D. Roux, G. S. Smith, S. K. Sinha, P. Dimon, N. A. Clark, and A. M. Belocq, Phys. Rev. Lett. **57**, 2718 (1986).

²W. Helfrich, Z. Naturforsch. Teil A **33**, 305 (1978).

³R. Lipowsky and S. Leibler, Phys. Rev. Lett. **56**, 2541 (1986).

⁴S. Leibler and R. Lipowsky, Phys. Rev. B (to be published).

⁵This has been confirmed in recent experiments, see D. Roux, in Proceedings of Les Houches Conference on Amphiphilic Films, Les Houches, France, February, 1987 (to be published).